

Solutions for Math 46 homework  
problems Day 11

Exercise 13 part b page 245

$$u(x) = b + \lambda \int_0^1 (u(y))^2 dy$$

The right hand side does not depend on  $x$ , so the left hand side also does not depend on  $x \Rightarrow u(x) \equiv A \leftarrow$  some constant

Now if  $\lambda = 0$  then we get

$A = b + 0 \Rightarrow u(x) \equiv b$  is the solution

If  $\lambda \neq 0$  we get

$$A = b + \lambda \int_0^1 A^2 dy = b + \lambda A^2$$

$$\lambda A^2 - A + b = 0$$

$$A_{1,2} = \frac{1 \pm \sqrt{1 - 4\lambda b}}{2\lambda} \quad (*)$$

Thus if  $\lambda \neq 0$  is such that  $1 - 4\lambda b > 0$  then we have two solutions  $u(x) = A_{1,2}$  given by  $x$   
this is

if  $1 = 4\lambda b$  then we have

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One solution  $u(x) = \frac{1}{2\lambda}$

if  $1 - 4\lambda b < 0$  then we have

no real valued solution

functions  $u(x)$



### Exercise 15 page 245

Solve the Fredholm equation

$$\int_0^1 \kappa(x, y) u(y) dy - \lambda u(x) = x$$

using eigenfunction expansions

where

$$\kappa(x, y) = \begin{cases} x(1-y) & \text{if } x < y \\ y(1-x) & \text{if } x > y \end{cases}$$

Solution

The kernel is symmetric and real valued. It is continuous since the only possible set of discontinuity points is along the line  $y = x$  where  $x(1-y)$  and  $y(1-x)$  match and

become  $x(1-x)$ .

Thus we can use Hilbert-Schmidt Theorem.

We search for eigen functions and eigen values

$$Ku = \int_0^1 k(x,y)u(y)dy$$

$$Ku = \lambda u$$

$$Ku(x) = \int_0^x y(1-x)u(y)dy + \int_x^1 x(1-y)u(y)dy$$

So  $(Ku(x))' \stackrel{\text{Leibnitz rule}}{=} \underbrace{x(1-x)u(x)}_{\text{cancel}} \frac{dx}{dx} - \underbrace{0(1-x)u(0)}_{=0} \frac{d0}{dx} + \int_0^x \frac{d}{dx} (y(1-x)u(y)) dy + \underbrace{x(1-1)u(1)}_{=0} \frac{d1}{dx} - \underbrace{x(1-x)u(x)}_{=} \frac{dx}{dx} + \int_x^1 \frac{d}{dx} (x(1-y)u(y)) dy =$

$-y'u(y)$

$-(1-y)u(y)$

$$= \int_0^x -y u(y) dy + \int_x^1 (1-y) u(y) dy$$

$$\text{Thus } (Ku(x))' = \int_0^x -y u(y) dy + \int_x^1 (1-y) u(y) dy$$

$$(Ku(x))'' = -xu(x) - (1-x)u(x) = -u(x)$$

$$\text{Thus } (Ku(x))'' = -u(x)$$

$$\text{but } Ku = \lambda u \Rightarrow (Ku)'' = \lambda u''$$

Thus we get  $\lambda u''(x) = -u(x)$   
 If  $\lambda = 0$  it has no nontrivial solutions  
 $Ku(0) = \lambda u(0)$  so wlog  $\lambda \neq 0$

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$$\int_0^0 y(1-0)u(y) dy + \int_0^1 0(1-y)u(y) dy = 0$$

$$Ku(1) = \lambda u(1)$$

$$\int_0^1 y(1-1)u(y) dy + \int_1^1 1(1-y)u(y) dy = 0$$

$$\Rightarrow u(1) = 0$$

Thus we get a Sturm  
Liouville problem

$$\left. \begin{aligned} u'' + \lambda u &= 0 \\ u(0) &= 0 \\ u(1) &= 0 \end{aligned} \right\}$$

As we already know the  
eigenvalues are  $\lambda = \left(\frac{\pi n}{1}\right)^2$   $n=1, 2, 3, \dots$

The eigenfunctions are

$$u_n(x) = \sin \frac{\pi n x}{1}$$

These are not of norm 1

$$(u_n, u_n) = \int_0^1 \sin^2(\pi n x) dx = \int_0^1 \frac{1}{2} - \frac{\cos(2\pi n x)}{2} dx$$

$$1 - 2 \sin^2 \alpha = \cos 2\alpha$$

$$\sin^2 \alpha = \frac{1 - \cos 2\alpha}{2}$$

$$\parallel \frac{1}{2}$$

$$\parallel u_n \parallel = \frac{1}{\sqrt{2}}$$

Note that Hilbert-Schmidt  
Theorem method works even  
when eigenfunctions are not of  
norm 1

Thus for the Fredholm equation we have

$$\lambda_n = \frac{1}{\lambda} = \frac{1}{(\pi n)^2} \leftarrow \text{eigenvalues.}$$

$$u_n(x) = \sin(\pi n x) \leftarrow \text{eigenfunctions}$$

$$n=1, 2, 3, 4, \dots$$

Let us find the presentation

$$x = \sum_{n=1}^{\infty} f_n u_n(x)$$

As we know from Fourier series discussions

$$f_n = \frac{(x, u_n)}{(u_n, u_n)} = 2(x, u_n)$$

$$(f, u_n) = \int_0^1 x \sin(\pi n x) dx = \frac{1}{2}$$

$(x, u_n)$

$$\left( -\frac{1}{\pi n} \cos(\pi n x) \right)'$$

$$= -\frac{x}{\pi n} \cos(\pi n x) \Big|_{x=0}^{x=1} + \int_0^1 \frac{\cos(\pi n x)}{\pi n} (x)' dx$$

$$= -\frac{1}{\pi n} \cos(\pi n) + \frac{\sin(\pi n x)}{(\pi n)^2} \Big|_{x=0}^{x=1} =$$

$$= \frac{(-1)^{n+1}}{\pi n}$$

"0"

Assume  $u(x) = \sum_{n=1}^{\infty} \alpha_n u_n(x)$

$$K(u(x)) = \lambda \left( \sum_{n=1}^{\infty} \alpha_n u_n(x) \right)$$

$$K\left(\sum_{n=1}^{\infty} \alpha_n u_n(x)\right) = \lambda \left(\sum_{n=1}^{\infty} \alpha_n u_n(x)\right)$$

$$\sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{\pi n} u_n(x)$$

$$\Rightarrow \sum_{n=1}^{\infty} \alpha_n K(u_n(x)) = \sum_{n=1}^{\infty} \alpha_n \lambda u_n(x) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{\pi n} u_n(x)$$

$$\sum_{n=1}^{\infty} \alpha_n \frac{1}{(\pi n)^2} u_n(x) = \sum_{n=1}^{\infty} \alpha_n \lambda u_n(x) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{\pi n} u_n(x)$$

$$\sum_{n=1}^{\infty} \left( \frac{\alpha_n}{(\pi n)^2} - \alpha_n \lambda \right) u_n(x) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{\pi n} u_n(x)$$

Compare the coefficients in front of  $u_n(x)$

$$\alpha_n \left( \frac{1}{(\pi n)^2} - \lambda \right) = \frac{(-1)^{n+1}}{\pi n}$$

$$\alpha_n = \frac{(-1)^{n+1}}{\frac{1}{(\pi n)^2} - \lambda}$$

(pages)

if  $\lambda \neq \frac{1}{(\pi n)^2}$

Thus if  $\lambda \neq \frac{1}{(\pi n)^2} \forall n$  then

the solution is

$$u(x) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{\frac{1}{(\pi n)^2} - \lambda} \sin(\pi n x)$$

if  $\lambda = \frac{1}{(\pi n)^2}$  then there is

no solution since the corresponding coefficient for the function  $x$

is  $\frac{(-1)^{n+1}}{\pi n} \neq 0$





Exercise 18 page 246

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Solve the integral equation

$$\int_0^1 e^{x+y} u(y) dy - \lambda u(x) = f(x)$$

by considering all cases

$$e^{x+y} = e^x e^y$$

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 $\alpha_1(x)$                  $\beta_1(y)$

So this is a

Fredholm equation with separable kernel. The matrix  $A$  is  $1 \times 1$

$$(A, \alpha_1) = \int_0^1 e^x e^x dx = \frac{1}{2} e^{2x} \Big|_{x=0}^{x=1} = \frac{1}{2} (e^2 - 1)$$

The column vector  $F$  is  $(\beta_1, f) =$

$$= \int_0^1 e^x f(x) dx$$

Thus equation 4.31 becomes

$$\sum_{j=1}^n \alpha_j(x) (A - \lambda I) \vec{c} = \vec{F}$$

$$\left( \frac{1}{2} (e^2 - 1) - \lambda \right) c = \int_0^1 e^x f(x) dx$$

The eigenvalue of  $A$  clearly is  $\tilde{\lambda} = \frac{1}{2}(e^2 - 1)$  (page 10)

Assume that  $\lambda \neq \tilde{\lambda}$  and  $\lambda \neq 0$

Then, as we know

$$c = \frac{\int_0^1 e^x f(x) dx}{\frac{1}{2}(e^2 - 1) - \lambda} \quad \text{and}$$

The solution is

$$\begin{aligned} u(x) &= \frac{1}{\lambda} (-f(x) + c \alpha_1(x)) = \\ &= \frac{1}{\lambda} \left( -f(x) + \frac{\int_0^1 e^y f(y) dy}{\frac{1}{2}(e^2 - 1) - \lambda} e^x \right) \end{aligned}$$

If  $\tilde{\lambda} = \frac{1}{2}(e^2 - 1)$  then the solution exists only if  $\vec{F}$  is in the range of  $\underbrace{A - \tilde{\lambda} I}_{\text{zero operator}} = 0$

i.e. only if  $\vec{F} = 0$  i.e. only if

$$\int_0^1 e^x f(x) dx = 0 \quad \text{in this case}$$

it is  $u(x) = \frac{1}{\tilde{\lambda}} (-f(x) + \tilde{c} e^x)$   
any constant

If  $\lambda = 0$  then the equation becomes

$$\int_0^1 e^x e^y u(y) dy - 0u(x) = f(x)$$

$$e^x \int_0^1 e^y u(y) dy = f(x)$$

↑ some constant

So the solution could exist only for  $f(x) = De^x$ . To find a solution in this case take any  $\tilde{u}(y)$  s.t.  $\int_0^1 e^y \tilde{u}(y) dy \neq 0$  multiply  $\tilde{u}(y)$  by a constant to make this integral  $D$ . This would be a particular solution. Plus you can add any solution to a homogeneous problem

$$e^x \int_0^1 e^y u(y) dy = 0$$

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Thus you can look at and  $u = \tilde{u} + u$  where  $u$  satisfies  $\int_0^1 e^y u(y) dy = 0$  😊