# LECTURE OUTLINE <br> Kinematics of Rotation 

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Math 15
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Goals

## Cross Product and the Determinant

Rigid Body Kinematics

## Last time we explored the Cross Product

Nice and Linear:

$$
\left(\vec{v}_{1}+c \vec{v}_{2}\right) \times \vec{w}=\vec{v}_{1} \times \vec{w}+c \vec{v}_{2} \times \vec{w}
$$

but, Not Commutative!

$$
\vec{v} \times \vec{w}=-\vec{w} \times \vec{v}
$$

and Not Associative!

$$
(\vec{u} \times \vec{v}) \times \vec{w}=\vec{u} \times(\vec{v} \times \vec{w})+\vec{v} \times(\vec{w} \times \vec{u})
$$

In particular, $\vec{u} \times \vec{v} \times \vec{w}$ has no good meaning!

## Computing a $3 \times 3$ Determinant

$$
\begin{aligned}
& \vec{u} \times \vec{v}=\left|\begin{array}{ll}
u_{y} & u_{z} \\
v_{y} & v_{z}
\end{array}\right| \hat{i}-\left|\begin{array}{cc}
u_{x} & u_{z} \\
v_{x} & v_{z}
\end{array}\right| \hat{j}+\left|\begin{array}{cc}
u_{x} & u_{y} \\
v_{x} & v_{y}
\end{array}\right| \hat{k} \\
& \vec{w} \cdot(\vec{u} \times \vec{v}) \equiv\left|\begin{array}{lll}
w_{x} & w_{y} & w_{z} \\
u_{x} & u_{y} & u_{z} \\
v_{x} & v_{y} & v_{z}
\end{array}\right|=\vec{w} \cdot\left|\begin{array}{ccc}
\hat{i} & \hat{j} & \hat{k} \\
u_{x} & u_{y} & u_{z} \\
v_{x} & v_{y} & v_{z}
\end{array}\right|
\end{aligned}
$$

Let $\vec{u}=\hat{i}-2 \hat{k}, \vec{v}=\hat{i}+\hat{j}+\hat{k}$ and $\vec{w}=3 \hat{i}+2 \hat{j}$, and compute $\vec{w} \cdot(\vec{u} \times \vec{v})$

## Application to spacial reasoning

Ex: Note we can think of a line as $t \vec{v}+\vec{p}$. Define what it means for two lines to be non-parallel, and find a formula for the distance between two non-parallel lines.

Find the distance between the lines $(\hat{i}-2 \hat{k}) t+\hat{i}$ and $\vec{v}=$ $t(\hat{i}+\hat{j}+\hat{k})$.

## Kinematics

For our rotation use $\hat{i}, \hat{j}$, and $\hat{k}$ with $\hat{k}$ the axis of rotation and $\hat{i}$ to $\hat{j}$ in the direction of rotation (right hand rule). Give $\hat{r}$ and $\hat{\theta}$ there usual meanings with respect to $\hat{i}, \hat{j}, \hat{k}$. There is a second "view" we might take to think about the center (center of mass) which we might call $\hat{i}_{c}, \hat{j}_{c}, \hat{k}_{c}$. A particle moving about another particle can be described by:

$$
\vec{r}=\vec{c}+r \hat{r}+z \hat{k} .
$$

## Kinematics

Usual Book Assumptions:
(1) $\dot{r}=0$ (Not Nutty)
(2) $z=0$ (Not so Nutty)
(3) $\frac{d \hat{k}}{d t}=0$ (Nutty)

Note: (3) The Great Yo-Yo Restriction allows us to assume $\frac{d \hat{i}}{d t}=0$ and $\frac{d \hat{j}}{d t}=0$ as well.

## Kinematics

## Book's Notation under the Usual Book Assumptions:

$$
\begin{gathered}
\vec{\omega} \equiv \dot{\theta} \hat{k} \equiv \omega \hat{k} \\
\vec{\alpha} \equiv \frac{d \vec{\omega}}{d t}=\ddot{\theta} \hat{k} \\
\vec{v} \equiv \frac{d \vec{r}}{d t}=r \dot{\theta} \hat{\theta}=\vec{\omega} \times \vec{r} \\
\vec{a} \equiv \frac{d \vec{v}}{d t}=-r \dot{\theta}^{2} \hat{r}+r \ddot{\theta} \hat{\theta}=\vec{\omega} \times \vec{v}+\vec{\alpha} \times \vec{r}
\end{gathered}
$$

## Derivatives in Polar Coordinates

$$
\begin{gathered}
\frac{d \hat{r}}{d t}=\dot{\theta} \hat{\theta} \\
\frac{d \hat{\theta}}{d t}=-\dot{\theta} \hat{r} \\
\frac{d \vec{r}}{d t}=\dot{r} \hat{r}+r \dot{\theta} \hat{\theta} \\
\frac{d^{2} \vec{r}}{d t^{2}}=\left(\ddot{r}-r \dot{\theta}^{2}\right) \hat{r}+(2 \dot{r} \dot{\theta}+r \ddot{\theta}) \hat{\theta}
\end{gathered}
$$

